# FAILURE OF AMALGAMATION IN HILBERT LATTICES 

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We show that Bruns and Harding's counterexample (see [1]) to amalgamation in orthomodular lattices also works for Hilbert lattices. The argument is based on the example of a non-modular Hilbert lattice devised by von Neumann in a letter to Birkhoff (see [2] for an extensive quotation from that letter). Consider the real sequence space $\ell_{2}=\left\{f: \mathbb{N}^{+} \rightarrow \mathbb{R}: \Sigma_{i=1}^{\infty} f(i)^{2}<\infty\right\}$, where $\mathbb{N}^{+}$stands for $\mathbb{N} \backslash\{0\}$. Let $\left\langle e_{n}: n \in \mathbb{N}^{+}\right\rangle$be the standard orthonormal base of $\ell_{2}$, i.e., $e_{n}(n)=1$ and $e_{n}(m)=0$ for $m \neq n$. It follows from Satz 15 in [3] that there are two unbounded self-adjoint operators $X$ and $Y$ such that $\operatorname{dom}(X) \cap \operatorname{dom}(Y)=\{0\}$ and moreover $Y$ can be chosen to be the "multiplication by $n$ " operator $Y f=\left\langle n \cdot f(n): n \in \mathbb{N}^{+}\right\rangle$. Thus, $Y$ is represented over the standard base as the $\mathbb{N}^{+} \times \mathbb{N}^{+}$matrix with entries $y_{i j}=i \cdot e_{i}(j)$. Further, the operators $X^{2}+2 I$ and $Y^{2}+2 I$ are self-adjoint and invertible. Define $A=\left(X^{2}+2 I\right)^{-1}$ and $B=\left(Y^{2}+2 I\right)^{-1}$, where $I$ is the identity operator. The following lemma spells out some properties of $A, B$ and $C$.
Lemma 1. The operators $A$ and $B$ are bounded self-adjoint, with $\|A\| \leq \frac{1}{2}$ and $\|B\| \leq \frac{1}{2}$. Further, $C=\sqrt{I-A^{2}-B^{2}}$ is also bounded and self-adjoint. Moreover the following conditions hold:
(1) $A^{2}+B^{2}+C^{2}=I$,
(2) $A^{-1}$ and $B^{-1}$ exist,
(3) $A f+B g=0$ implies $f=g=0$,
(4) $C f=0$ implies $f=0$,
(5) if $B^{-1}$ is defined for $f$ and $B^{-1} f=0$, then $f=0$,
(6) if $B^{-1}$ is defined for $f$ and $C B^{-1} f=0$, then $f=0$.

Let now $\bar{H}$ be the direct sum $\ell_{2} \oplus \ell_{2} \oplus \ell_{2}$. From (1) it follows that the map $f \mapsto(A f, B f, C f)$ is an embedding of $\ell_{2}$ into $\bar{H}$. Let $\bar{L}$ be the image of $\ell_{2}$ under this embedding. Clearly, $\bar{L}$ is a proper subspace of $\bar{H}$. Let $\bar{K}=\left\{(0,0, h): h \in \ell_{2}\right\}$ and $\bar{M}=\left\{(0, g, h): g, h \in \ell_{2}\right\}$.

Lemma 2. The spaces $\bar{K}, \bar{L}$ and $\bar{M}$ satisfy the following conditions:

- $\bar{K} \subset \bar{M}$,
- $\bar{L} \vee \bar{K}=\overline{1}$,
- $\bar{L} \cap \bar{M}=\overline{0}$.
where $\overline{1}$ stands for $\bar{H}$ and $\overline{0}$ for the zero subspace.
In other words, $\overline{1}, \bar{K}, \bar{L}, \bar{M}$, and $\overline{0}$ form a sublattice of the Hilbert lattice of $\bar{H}$ isomorphic to the pentagon.

To proceed, we will need certain derived operators, in particular $C B^{-2} C$. Since $C$ in bounded self-adjoint, it is represented over the standard base by a Hermitian
$\mathbb{N}^{+} \times \mathbb{N}^{+}$matrix, say

$$
C=\left(\begin{array}{ccc}
c_{11} & c_{12} & \ldots \\
c_{21} & c_{22} & \ldots \\
\vdots & \vdots & \ddots
\end{array}\right)
$$

with $c_{i j}=c_{j i}$. The operator $B^{-2}=B^{-1} B^{-1}=\left(Y^{2}+2 I\right)\left(Y^{2}+2 I\right)$ can also be represented by a Hermitian (unbounded, diagonal) matrix

$$
\left(\begin{array}{cccc}
9 & 0 & 0 & \ldots \\
0 & 36 & 0 & \ldots \\
0 & 0 & 121 & \ldots \\
\vdots & \vdots & \vdots & \ddots
\end{array}\right)
$$

with main diagonal entries $p_{i i}=\left(i^{2}+2\right)^{2}$. Therefore, the operator $C B^{-2} C$ is represented by the (non-Hermitian, unbounded) matrix

$$
M=\left(\begin{array}{ccc}
m_{11} & m_{12} & \ldots \\
m_{21} & m_{22} & \ldots \\
\vdots & \vdots & \ddots
\end{array}\right)
$$

where $m_{i j}=\sum_{n=1}^{\infty}\left(n^{2}+2\right)^{2} c_{i n} c_{n j}$.
Lemma 3. The space $\bar{L}^{\perp} \cap \bar{M}$ is a non-zero subspace of $\bar{H}$.
Let $\bar{N}$ be $\bar{L}^{\perp} \cap \bar{M}$. Since $\bar{N}$ is a non-zero proper subspace of $\bar{M}$, by properties of orthomodular lattices it follows that $\bar{N}$ and $\bar{M}$ generate an 8-element Boolean subalgebra of (the Hilbert lattice of) $\bar{H}$. In particular, putting $\bar{U}=\bar{N} \vee \bar{M}^{\perp}$ we get that $\bar{U}$ is a proper subspace of $\bar{H}$.

Lemma 4. The space $\bar{U}$ is precisely $\{(a, b, c): B b+C c=0\}$, with $a, b, c \in \ell_{2}$. Moreover, the following equalities hold:
(1) $\bar{K} \cap \bar{U}=0$,
(2) $\bar{L} \cap \bar{U}=0$,
(3) $\bar{U} \cap \bar{K}^{\perp}=\bar{M}^{\perp}$,
(4) $\bar{U}^{\perp} \vee \bar{L}=\bar{N}^{\perp}$.

We now preceed to define a subspace of $\bar{U}$. Let $P$ be the projection on first coordinate, i.e., an operator on $\ell_{2}$ defined by

$$
\operatorname{Pf}(n)= \begin{cases}f(n) & \text { if } n=1 \\ 0 & \text { otherwise }\end{cases}
$$

Clearly the subspace $\bar{P}$ corresponding to $P$ is a proper, non-zero subspace of $\ell_{2}$. Thus, $\bar{W}$ defined as $\{(a, b, P c): B b+C c=0\}$ is a proper, non-zero subspace of $\bar{U}$.

Lemma 5. We have $(\bar{L}+\bar{W})^{\perp}=0=(\bar{K}+\bar{W})^{\perp}$.
Lemma 6. The subalgebra of the Hilbert lattice of $\bar{H}$ generated by $\bar{L}, \bar{K}$ and $\bar{W}$ is isomorphic to MO3 (i.e., the Chinese lantern with 6 atoms).

Now let $\mathbf{L}_{1}$ and $\mathbf{L}_{2}$ be two copies of the Hilbert lattice of $\bar{H}$. We will use subscripts to differentiate between their elements (subspaces of the two copies of
$\bar{H})$. Thus, for instance $\bar{L}_{1}$ will stand for $\bar{L} \subset \bar{H}$ as an element of $\mathbf{L}_{2}$, while $\bar{L}_{2}$ will stand for $\bar{L}$ as an element of $\mathbf{L}_{2}$.

By Lemma 6, MO3 is a sublattice of both. Let $p, q, r$ be three distinct atoms of MO3. Define embeddings $e_{1}: M O 3 \longrightarrow \mathbf{L}_{1}$ and $e_{2}: M O 3 \longrightarrow \mathbf{L}_{2}$ by putting:

- $e_{1}(p)=\bar{L}_{1}, e_{1}(q)=\bar{K}_{1}, e_{1}(r)=\bar{W}_{1}$, and
- $e_{2}(p)=\bar{L}_{2}^{\perp}, e_{2}(q)=\bar{K}_{2}^{\perp}, e_{1}(r)=\bar{W}_{2}^{\perp}$.

Notice that intuitively this amounts to identifying the copies of $\bar{L}, \bar{K}$ and $\bar{W}$ in $\mathbf{L}_{1}$ respectively with the copies of $\bar{L}^{\perp}, \bar{K}^{\perp}$ and $\bar{W}^{\perp}$ in $\mathbf{L}_{2}$.

Suppose that the V-formation $\left\langle M O 3, \mathbf{L}_{1}, \mathbf{L}_{2}, e_{1}, e_{2}\right\rangle$ can be amalgamated by a Hilbert lattice $\mathbf{L}$. We will use $\overline{1}$ and $\overline{0}$ respectively for the top and bottom elements of our lattices. In $\mathbf{L}$ we can then carry out the following calculations.

$$
\begin{equation*}
\overline{1}=\bar{U}_{2}^{\perp} \vee \bar{K}_{2}^{\perp}=\bar{U}_{2}^{\perp} \vee q \leq \bar{U}_{2}^{\perp} \vee q \vee \bar{U}_{1}^{\perp}=\bar{U}_{2}^{\perp} \vee \bar{K}_{1} \vee \bar{U}_{1}^{\perp}=\bar{U}_{2}^{\perp} \vee \bar{M}_{1} \tag{1}
\end{equation*}
$$

where the inequality in the middle is trivial, the equalities neighbouring it follow from the embeddings, and the first and last equalities in the row both come from Lemma 4 (from (1) and (3) respectively). Since $\bar{N}_{1} \subseteq \bar{M}_{1}$, by orthomodularity we get

$$
\begin{equation*}
\bar{N}_{1} \vee \bar{U}_{1}^{\perp}=\bar{N}_{1} \vee\left(\bar{N}_{1}^{\perp} \cap \bar{M}_{1}\right)=\bar{M}_{1} \tag{2}
\end{equation*}
$$

which together with (1) yields

$$
\begin{equation*}
\overline{1}=\bar{U}_{2}^{\perp} \vee \bar{N}_{1} \vee \bar{U}_{1}^{\perp} \tag{3}
\end{equation*}
$$

Then, as $\bar{N}_{1} \subseteq \bar{U}_{1}$ and $\bar{U}_{2}^{\perp} \subseteq \bar{W}_{2}^{\perp}=r=\bar{W}_{1} \subseteq \bar{U}_{1}$, using orthomodularity again we obtain

$$
\begin{equation*}
\bar{N}_{1} \vee \bar{U}_{2}^{\perp}=\bar{U}_{1} . \tag{4}
\end{equation*}
$$

Orthomodularity also yields

$$
\begin{equation*}
\bar{L}_{1}^{\perp}=\bar{N}_{1} \vee\left(\bar{L}_{1}^{\perp} \cap \bar{N}_{1}^{\perp}\right) . \tag{5}
\end{equation*}
$$

Therefore

$$
\begin{aligned}
\bar{N}_{2}^{\perp}=\bar{U}_{2}^{\perp} \vee \bar{L}_{2} & =\bar{U}_{2}^{\perp} \vee p^{\perp}=\bar{U}_{2}^{\perp} \vee \bar{L}_{1}^{\perp} \\
=\bar{U}_{2}^{\perp} \vee\left(\bar{N}_{1} \vee\left(\bar{L}_{1}^{\perp} \cap \bar{N}_{1}^{\perp}\right)\right) & =\left(\bar{U}_{2}^{\perp} \vee \bar{N}_{1}\right) \vee\left(\bar{L}_{1}^{\perp} \cap \bar{N}_{1}^{\perp}\right) \\
=\bar{U}_{1} \vee\left(\bar{L}_{1}^{\perp} \cap \bar{N}_{1}^{\perp}\right) & =\left(\bar{M}_{1}^{\perp} \vee \bar{N}_{1}\right) \vee\left(\bar{L}_{1}^{\perp} \cap \bar{N}_{1}^{\perp}\right) \\
=\bar{M}_{1}^{\perp} \vee\left(\bar{N}_{1} \vee\left(\bar{L}_{1}^{\perp} \cap \bar{N}_{1}^{\perp}\right)\right) & =\bar{M}_{1}^{\perp} \vee \bar{L}_{1}^{\perp} \\
=\left(\bar{M}_{1} \cap \bar{L}_{1}\right)^{\perp} & =\overline{0}^{\perp}=\overline{1}
\end{aligned}
$$

where the first equality follows by Lemma $4(4)$ and the last row employs Lemma 2. So we obtain $\bar{N}_{2}^{\perp}=\overline{1}$, which contradicts Lemma 3.

## References

[1] G. Bruns, J. Harding, Amalgamation of ortholattices, Order, 14, 193-209, 1998.
[2] M. Redei, The birth of quantum logic, manuscript.
[3] J. von Neumann, Zur Theorie der unbeschränkter Matrizen, Collected Works, vol. II, no. 3, 144-172, Pergamon Press 1961.

