TRACING THE OUTER HALO IN A GIANT ELLIPTICAL TO $25 R_{\text{eff}}$

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ABSTRACT

We have used the Advanced Camera for Surveys and Wide Field Camera 3 cameras on board the *Hubble Space Telescope* to resolve stars in the halo of the nearest giant elliptical (gE) galaxy NGC 5128 out to a projected distance of 140 kpc (25 effective radii, R_{eff}) along the major axis and 90 kpc (16 R_{eff}) along the minor axis. This data set provides an unprecedented radial coverage of the stellar halo properties in any gE galaxy. Color–magnitude diagrams clearly reveal the presence of the red giant branch stars belonging to the halo of NGC 5128, even in our most distant fields. The star counts demonstrate increasing flattening of the outer halo, which is elongated along the major axis of the galaxy. The $V - I$ colors of the red giants enable us to measure the metallicity distribution in each field and so map the gradient out to $\sim 16 R_{\text{eff}}$ from the galaxy center along the major axis. A median metallicity is obtained even for the outermost fields along both axes. We observe a smooth transition from a metal-rich ($[M/H] \sim 0.0$) inner galaxy to lower metallicity in the outer halo, with the metallicity gradient slope along the major axis of $\Delta[M/H]/\Delta R \simeq -0.0054 \pm 0.0006 \text{ dex kpc}^{-1}$. In the outer halo, beyond $\sim 10 R_{\text{eff}}$, the number density profile follows a power law, but also significant field-to-field metallicity and star count variations are detected. The metal-rich component dominates in all observed fields, and the median metallicity is $[M/H] > -1 \text{ dex}$ in all fields.

Key words: galaxies: elliptical and lenticular, cD – galaxies: halos – galaxies: individual (NGC 5128) – galaxies: stellar content

Online-only material: color figures

1. INTRODUCTION

Because dynamical timescales in the outer halos of large galaxies are long, the properties of stars in the remote halo regions, such as their metallicity and surface density distributions and gradients, provide important constraints for galaxy formation and hierarchical assembly models (Johnston et al. 2008; Font et al. 2011; Naab & Ostriker 2009; Cooper et al. 2013). However, the extremely faint surface brightness of the halo makes the raw observations challenging. Integrated-light studies need to probe surface brightness levels of $\mu \sim 30\text{--}32 \text{ mag arcsec}^{-2}$, predicted in simulations for outer halos (e.g., Johnston et al. 2008) where most of the sub-structure tracing the hierarchical assembly history of the galaxies is expected. Such faint limits are out of reach for integrated light studies with most modern large telescopes (Jablonka et al. 2010; Abraham & van Dokkum 2014), though attempts have been made to derive *mean* properties by stacking many similar galaxies (Zibetti et al. 2004; Tal & van Dokkum 2011).

A powerful alternative to integrated light studies is resolved photometry of individual stars, a technique that can be applied for galaxies at distances $\lesssim 20 \text{ Mpc}$. Direct stellar photometry is capable of revealing not just mean halo properties but the *distribution functions* of metallicity and age.

Only a few giant elliptical (gE) galaxies are close enough to allow detection and study of their outer-halo regions. Mihos et al. (2013) traced the halo color gradient of the Virgo Cluster gE M49 out to $\sim 100 \text{ kpc}$ ($7 R_{\text{eff}}$), where $\mu_V \sim 28 \text{ mag arcsec}^{-2}$. The bluer color of the outer halo and the steepening of the color gradient in the M49 halo, if interpreted as purely due to metallicity, argues for the presence of a dominant metal-poor

population with $[\text{Fe}/\text{H}] < -1$. The detection of this “classical” metal-poor halo is in contrast with the findings in the Centaurus A (Cen A) group central giant NGC 5128, where both a high mean metallicity and no strong metallicity gradient have been seen out to a similar scaled distance of $6.5 R_{\text{eff}}$ (Rejkuba et al. 2005). In the Leo group gE NGC 3379 the transition to a classically metal-poor outer halo was found at $\sim 10\text{--}12 R_{\text{eff}}$ (Harris et al. 2007). These few galaxies and very limited spatial coverage (with only one or few pointings per galaxy) preclude drawing general conclusions and demand a wider range of observations.

NGC 5128 (often referred to as Cen A, after its radio source name) is a particularly important target. A mere 3.8 Mpc distant (Harris et al. 2010), it is by far the closest and easiest to observe gE galaxy. Its old-halo red-giant stars can be studied readily with *Hubble Space Telescope* (HST) imaging (e.g., Soria et al. 1996; Marleau et al. 2000; Harris et al. 1999) or under excellent seeing conditions from the ground (Rejkuba et al. 2001, 2003; Crnojević et al. 2013). HST-based photometry, reaching at least 1–1.5 mag below the tip of the red giant branch (RGB), has allowed the metallicity distribution function (MDF) to be investigated in detail in four fields at projected galactocentric distances from 8 to 38 kpc (Harris et al. 1999; Harris & Harris 2000, 2002; Rejkuba et al. 2005). The outer-halo stars were traced out to $\sim 85 \text{ kpc}$ with VIMOS at the ESO Very Large Telescope (Crnojević et al. 2013) and found to have high median metallicity even at $\sim 15 R_{\text{eff}}$.⁸ This VIMOS study also provides a first hint that the outer halo becomes more elongated and has a

⁸ Crnojević et al. (2013) used $R_{\text{eff}} = 330'' = 6.1 \text{ kpc}$, while we adopt $305'' = 5.6 \text{ kpc}$ from Dufour et al. (1979).

Figure 1. Distribution of the *HST* observations in the halo of NGC 5128 relative to the center of the galaxy projected on the plane of the sky. The elliptical contours have an axis ratio of 0.77 and P.A. = 35° and are plotted for 1.5, 4.0, 5.5, 7.0, 10.5, 15.5, and 25 R_{eff} distance, which are representative for our present and past *HST* pointings. The two green stars are dwarf spheroidal galaxies in the Cen A group, KK197 and KK55.

(A color version of this figure is available in the online journal.)

surface brightness profile shallower than the $r^{1/4}$ law. However, high contamination by both foreground stars and unresolved background galaxies prevented stronger conclusions. Deeper and higher-resolution *HST* observations can go much further.

In this work, we extend our previous *HST* studies of the NGC 5128 halo into the extreme outer parts of the halo up to 140 kpc or 25 R_{eff} , regions never observed before. With these new data we address the following questions: (1) How far does the halo extend? (2) Is there a genuine metal-poor halo in NGC 5128, similar to that found in the outer regions of the Milky Way (MW), M31, NGC 3379, and M49 (Carollo et al. 2007; Chapman et al. 2006; Harris et al. 2007; Mihos et al. 2013)? (3) If so, at which radius does the metal-poor population start to dominate?

2. OBSERVATIONS AND DATA REDUCTION

Five new fields were observed with the *HST* during Cycle 20 (Cy20; GO program 12964) with the Wide Field Camera 3 (WFC3) in UVIS mode as primary and Advanced Camera for Surveys (ACS) in WFC mode as the secondary (parallel) instrument. For both instruments imaging was done with the F606W and F814W, the same filters as used in all our previous studies. Each field received three dithered exposures adding to one orbit per filter. The total exposure time for WFC3 was 2376 s (F814W) and 2496 s (F606W), while the ACS images got 2137 and 2270 s in the same two bands.

Figure 1 shows the locations of our fields relative to the center of NGC 5128 projected on the plane of the sky. Fields F1–F4 are described in Harris et al. (1999), Harris & Harris (2000, 2002), and Rejkuba et al. (2005), while our new Cy20 fields F5–F9 are marked with suffix A (ACS) or W (WFC3). Following Dufour et al. (1979), the contour lines shown are for an axis ratio 0.77, with the major axis along a position angle of 35°. The effective

radius is $R_{\text{eff}} = 305''$ corresponding to 5.6 kpc at the 3.8 Mpc distance (Harris et al. 2010). Our new fields F5–F7 are located along the major axis ~ 60 , 90, and 140 kpc northeast from the center of NGC 5128, respectively. The fields F8 and F9 along the northwest minor axis lie on the same ellipses as F4 and F6, which are located respectively at ~ 40 ($\sim 7 R_{\text{eff}}$) and 90 kpc ($\sim 16 R_{\text{eff}}$) close to and along the major axis.

The pipeline processed images were downloaded from the *HST* archive, and pixel-based charge transfer corrections (Anderson & Bedin 2010) were applied with the Fortran code⁹ from J. Anderson (Ubeda & Anderson 2012). Point-spread-function (PSF)-fitting photometry was completed with the DAOPHOT suite of programs (Stetson 1987). Our images are completely uncrowded, so in principle aperture photometry would be accurate, but PSF/ALLSTAR fitting was preferred to allow us using the DAOPHOT shape parameters (χ , sharp) to cull nonstellar objects (faint, small background galaxies) as far as possible. Stars detected were retained if (1) their positions on both filters matched to within 1 pixel, (2) $\chi < 1.5$, ABS (sharp) < 0.4 , and (3) measurement uncertainties were smaller than 0.5 mag. Photometric calibration in the VEGAMAG system followed the prescription in Sirianni et al. (2005), with the latest zero points recommended on the instrument Web pages.¹⁰ Final transformation to the ground-based *VI* magnitude scale follows Saha et al. (2011).

3. RESULTS

3.1. Color–Magnitude Diagrams

In Figure 2 we show the color–magnitude diagrams (CMDs) for the new fields, in comparison with the deep CMD in F4 from Cycle 11 (Rejkuba et al. 2005). The top row displays the CMDs in fields along the major axis (F5–F7), while the bottom row shows the CMDs for the two fields along the minor axis (F8–F9). The CMDs for fields along the minor axis fields are plotted immediately below the corresponding CMD of the major axis field that is on the same ellipse (see Figure 1). The magnitudes and colors are extinction-corrected for each field individually following the Schlafly & Finkbeiner (2011) recalibration of the Galactic extinction maps of Schlegel et al. (1998), assuming the Fitzpatrick (1999) reddening law with $R_V = 3.1$. The dotted horizontal line marks the expected magnitude of the tip of the RGB at $I_0 = 23.88$ (Harris et al. 2010).

The CMDs in Figure 2 show a very similar, broad upper RGB in fields F4 and F8; F5 is almost as broad in color. F6, however, has a narrower upper RGB, in particular in the ACS pointings (black dots), suggesting the beginning of a transition to a more metal-poor-dominated population. In F7 and F9, the RGB is quite sparsely populated, but the color distribution still resembles that of F6.

Most of the objects brighter than $I_0 = 23.88$ (i.e., clearly above the RGB tip) are foreground MW stars, but a few bright asymptotic giant branch (AGB) stars are expected in NGC 5128 halo, based on the study of long-period variables (Rejkuba et al. 2003), and 38 kpc field star formation history (Rejkuba et al. 2011). The camera areas ratio (ACS/WFC3) is 1.55, and we thus expect the same ratio for the number of bright stars, if they all belong to the MW foreground. This is indeed observed in F7 and F9, but there are $1.8\times$ fewer stars in F5, F6, and F8

⁹ http://www.stsci.edu/hst/wfc3/tools/cte_tools

¹⁰ http://www.stsci.edu/hst/wfc3/phot_zp_lbn;
<http://www.stsci.edu/hst/acs/analysis/zeropoints>

Figure 2. Color–magnitude diagrams of the stellar sources detected in the fields in the halo of NGC 5128. F4 is the ACS field previously obtained by Rejkuba et al. (2005), while the other five panels show the new CMDs observed with WFC3 (blue open circles) and ACS (black dots) during our Cy20 *HST* program. The dotted horizontal line indicates the tip of the RGB.

(A color version of this figure is available in the online journal.)

WFC3 pointings compared to ACS. This might be due to a small variation of contribution from bright AGB stars in NGC 5128 or sub-structure in the MW foreground, and will be investigated in a follow-up study. At levels fainter than the tip of the RGB ($I_0 > 23.88$), foreground stars are also present, and at the faintest magnitudes ($I \gtrsim 27$) compact high-redshift galaxies may also add to the contamination.

To estimate the foreground contamination quantitatively, we use the TRILEGAL (Girardi et al. 2005) and Besançon (Robin et al. 2003) MW models, while the contamination by background unresolved galaxies can be evaluated from the Hubble Ultra Deep Field (UDF) observations.¹¹ We transformed the UDF F775W AB measurements to F814W on Vegamag photometric system adopting the relation from Sirianni et al. (2005). An example of the model field components is shown in Figure 3 for F9. The model star counts above the RGB tip reveal significant differences between TRILEGAL and Besançon. Though the investigation of these differences is beyond the scope of this work, we decided to adopt TRILEGAL (version 1.6), because it had about 30% more stars with respect to the Besançon simulation, and it better matched the bright end of the observed luminosity function (Figure 3, left), and the color distribution of the bright foreground stars with $20 \leq I < 22$ (Figure 3, right).

The model field population from TRILEGAL run at each field location, normalized to the ACS/WFC and WFC3 areas, was used to statistically clean the observed CMDs in all five fields. This step removed almost all stars brighter than the RGB tip, except stars redder than $V - I > 2.8$, which are missing in the TRILEGAL simulations. The simulations include reddening, but were not corrected for incompleteness and therefore it is likely that the faintest stars are oversubtracted. However, since we are interested in the upper ~ 1.6 mag below the RGB tip, from which we measure the metallicity distributions, and the

photometry is highly complete at that level, any such corrections will have no significant effect on the MDFs. Remaining potential contaminants from unresolved galaxies are mostly fainter and bluer than the stars used to derive MDFs (Figure 3, right). This, combined with their low total number shows that the unresolved galaxies do not affect our metallicity distribution derivation.

The fact that the RGB is still detectable in F7 at 140 kpc ($25 R_{\text{eff}}$) is noteworthy by itself; our original expectation in designing the observations was that it would act as a “control field” of almost pure field population based on outward extrapolation of the inner-halo $r^{1/4}$ density profile. In short, we have not yet found the limits of this giant galaxy’s halo. It is also striking that the CMD in the F9 minor axis field, lying on the same ellipse as F6 (the major axis field at $16 R_{\text{eff}}$), appears similar to F7 at the much larger distance of $25 R_{\text{eff}}$ along the *major* axis field; if the isophotal contours kept the same ellipticity with radius, then the level of star counts in F9 should be the same as in F6, but instead it is considerably lower, and similar to that in F7 (see Section 3.2). This evidence suggests that either the outer halo becomes increasingly elongated, as already suspected from the VIMOS observations by Crnojević et al. (2013); or that some sort of very extended tidal debris plume lies along the major axis, boosting the level of star counts over the normal smooth halo component.

3.2. Halo Metallicity and Number Density Profiles

Our analysis of the very deep F4 photometry, which reached the red clump helium burning stars, quantitatively demonstrated that 70%–80% of the stars have classically “old” ages in excess of ~ 10 Gyr (Rejkuba et al. 2011). The colors of the old RGB stars can therefore be used also for other halo fields to derive the MDF without complications from the age/metallicity degeneracy. Conversion of $(V - I)_0$ to heavy-element abundance Z for stars with $-3.6 < M_{\text{bol}} < -2$ was done by interpolation in an isochrone grid as described in Harris & Harris (2000), with

¹¹ <http://heasarc.gsfc.nasa.gov/W3Browse/hst/hubbleudf.html>

Figure 3. Left: the observed I -band luminosity function in the F9 ACS field (black hashed histogram) compared with the simulated MW foreground luminosity function from the TRILEGAL v1.6 (blue), and Besançon model (red) of the Galaxy, and with the UDF point-like sources (green hashed histogram). The vertical dotted line at $I = 24.1$ indicates the position of the RGB tip magnitude in NGC 5128. Right: the color distribution of the F9 ACS field (black hashed histogram) in different magnitude bins compared with the simulated foreground stars from TRILEGAL v1.6 (blue), Besançon (red), and UDF point-like sources (green histogram). Simulated and observed distributions include reddening.

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Figure 4. Metallicity distribution functions derived from contamination-cleaned CMDs. The red (dashed) histogram in each panel is the F4 MDF. The blue is for the WFC3 data and black for ACS. For F7 and F9, the ACS and WFC3 measurements are combined (solid black histogram). All histograms are normalized by the total number of stars shown in each panel.

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the α -enhanced Teramo isochrones (Pietrinferni et al. 2006) as done for F4 (Rejkuba et al. 2011). The MDFs for the inner fields 1–3 (observed with WFPC2) have also been re-derived with this newer isochrone set. The absolute metallicity scale coming from the use of the new set of isochrones results in slightly higher average metallicity (by about 0.2 dex) than reported in Rejkuba et al. (2005). However, all metallicities derived in this work

are on the same scale and thus a *relative* comparison between the distributions of different fields is robust and can be used to investigate the halo metallicity gradient.

Figure 4 shows the new MDFs for all our fields F4–F9, derived from the foreground cleaned CMDs. The metallicity distributions are wide in all the panels, but in F7 and F9 they are poorly sampled because of small number statistics. We thus

Figure 5. Median metallicity (left) and density profile of RGB stars with $24 \leq I_0 < 26$ (right) as functions of the projected distance from the center of NGC 5128 along the major axis. Open green squares indicate the minor axis fields, reported at their projected distance multiplied by the inverse axis ratio (a/b), thus referred to the ellipse to which they belong (Figure 1). In the left panel, crosses show the results obtained by combining the ACS and WFC3 data for F7 and F9 fields. In the right panel we plot different fits to the density profile: a $R^{1/4}$ law applied to the inner 6 R_{eff} (solid blue), a power law applied to all (short dashed red) and only major axis (long dashed black) fields.

(A color version of this figure is available in the online journal.)

show combined WFC3 and ACS MDFs for F7 and F9. For comparison in all panels we overplot the MDF measured from ACS F4 data set.

The left panel of Figure 5 shows the *median* metallicity as a function of projected distance from the center of the galaxy along the major axis. The green open symbols denote minor axis fields (F8 and F9), for which the distance is scaled to the major axis by multiplying with the inverse axis ratio ($a/b = 1/0.77$). A clear metallicity gradient exists, with slope $\Delta[M/H]/\Delta R \simeq -0.0054 \pm 0.0006 \text{ dex kpc}^{-1}$, or $\Delta[M/H]/\Delta R_{\text{eff}} \simeq -0.030 \pm 0.003 \text{ dex per } R_{\text{eff}}$. Remarkably, however, even in the most remote parts of the halo we have sampled, the median metallicity has not fallen below $[M/H] = -1 \text{ dex}$.

Field-to-field scatter also shows up, such as much lower $[M/H]$ in the ACS versus WFC3 pointing of F6. For F7 and F9 fields we plot in addition to individual ACS and WFC3, also a median $[M/H]$ from both cameras together (crosses). Bootstrap simulations with 5000 random extractions of smaller samples representative for F7 and F9 MDFs, show that the median metallicity can be derived with $\lesssim 0.1 \text{ dex}$ accuracy for samples with ≥ 40 stars, while the uncertainty increases to $\sim 0.2 \text{ dex}$ for samples of 20–40 stars. These simulations demonstrate that the metallicity difference between F7 and F9 is significant in spite of low statistics.

The difference between the F7 and F9 mean metallicity is even more remarkable considering that these two fields have similar stellar density. The right panel of Figure 5 shows the number density of RGB stars with $24 \leq I_0 < 26$ (selected from foreground cleaned CMDs) as a function of the projected distance along the major axis. The previously mentioned striking difference between the minor and major axis fields located at the same projected major axis distance (F6 and F9) is even better

appreciated here. The minor axis fields density profile follows the extrapolation of $r^{1/4}$ profile, while a power law provides a better fit along the major axis beyond $R_{\text{eff}} \sim 10$, where the stellar density exceeds the shallower $r^{1/4}$ profile. Additional field-to-field variations seen in metallicity are also observed in density distribution gradient. We may, perhaps, be seeing the record of incompletely mixed debris from hierarchical buildup of the halo.

4. DISCUSSION AND CONCLUSIONS

The new *HST* observations presented here trace the halo of an early-type giant galaxy to unprecedentedly large distances ($25 R_{\text{eff}}$). It is interesting to put this distance in a context of cosmological “virial radius” R_{200} , at which the mean density equals $200\times$ the cosmological critical density. Taking the total mass of $(9.2 \pm 0.3) \times 10^{12} M_{\odot}$ for the entire Cen A group (Woodley 2006, using dynamics of all the satellites), would imply that the Cen A group virial radius is $R_{200} = 400 \text{ kpc}$, and our outermost field at 140 kpc corresponds to one-third of the virial radius.

Comparison of the RGB star counts between the major and minor axes suggests increasing elongation of the halo beyond $\sim 10 R_{\text{eff}}$. While sub-structure is expected to be present in the outer halo, the significantly lower density of stars at $16 R_{\text{eff}}$ along the minor versus major axis is also confirmed by the detection of stars at much larger distance along the major axis. This result furthermore confirms a tentative detection of the increasing ellipticity in the outer halo by Crnojević et al. (2013). This result is also interesting in comparison with the average ellipticity of stellar halos of early-type galaxies from Tal & van Dokkum (2011), which is based on a very deep stack of more than 42,000 Sloan Digital Sky Survey images of luminous

red galaxies. These authors found an increasing ellipticity with radius from $\epsilon = 0.25$ at 10 kpc to $\epsilon = 0.30$ at 100 kpc. The ellipticity in the halo of NGC 5128 appears to increase even more, although our so-far sparse sampling of the halo is only suggestive.

Tal & van Dokkum (2011) furthermore found that the surface brightness profile of their average stellar halo follows a Sersic model out to 100 kpc, but an excess of light beyond that is present in r -, i -, and z -band stacks. They discuss whether this excess light is part of the galaxy haloes or it traces the unresolved intragroup or intracluster light. With our measurement of the NGC 5128 metallicity gradient and star counts we find no clear transition to intragroup light beyond 90 kpc (beyond F6) in NGC 5128. Observations of further more distant fields along the minor and major axis would be necessary, however, to confirm this.

While the inner parts of the galaxy show a smooth metallicity gradient, we observe quite strong field-to-field variations in the median metallicity of the stars in the outer regions, as expected if the halo is built up by accretion of different satellites. However, the relatively high median metallicity is surprising, and the extended halo along the major axis might suggest that we are looking at least partly at material ejected from a long-ago accretion of a disk galaxy. This conclusion is supported by the analysis of the age distributions of stars (Rejkuba et al. 2011) and globular clusters (Woodley et al. 2010), as well as the presence of dynamical subgroupings of the globular clusters and planetary nebulae in the halo (Woodley & Harris 2011).

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Facilities: *HST* (ACS), *HST* (WFC3)

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