Low energy electron energy-loss spectroscopy of $\text{CF}_3X$ ($X=\text{Cl}, \text{Br}$)

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We report threshold electron energy-loss spectra for the fluorohalomethanes $\text{CF}_3X$ ($X=\text{Cl}, \text{Br}$). Measurements were made at incident electron energies of 30 and 100 eV in energy-loss range of 4–14 eV, and at scattering angles of 4° and 15°. Several new electronic transitions are observed which are ascribable to excitation of low-lying states as well as are intrinsically overlapped in the molecules themselves. Assignments of these electronic transitions are suggested. These assignments are based on present spectroscopic and cross-section measurements, high-energy scattering spectra, and ab initio molecular orbital calculations. The calculated potential curves along the C–X bond show repulsive nature, suggesting that these transitions may lead to dissociation of the C–X bond. The present results are also compared with the previous ones for $\text{CF}_3\text{H}$, $\text{CF}_4$, and $\text{CF}_3\text{I}$. © 2007 American Institute of Physics. [DOI: 10.1063/1.2424704]

I. INTRODUCTION

Freon and related halogen-substituted methane molecules $\text{CF}_3X$ ($X=\text{H}, \text{F}, \text{Cl}, \text{Br}, \text{I}$) are, in general, important industrial molecules with wide-ranging applications such as semiconductor etching, refrigerants, fire extinguishers, etc. These tetrahedrolike systems are of fundamental interest for studying halogen substitutional effects as well. However, it has been widely recognized that $\text{CF}_3\text{Cl}$ and $\text{CF}_3\text{Br}$ are responsible, via photochemical mechanisms, for the depletion of the earth’s ozone layer that protects our planet from the damaging effects of short wavelength solar radiation. Also $\text{CF}_3\text{H}$ and $\text{CF}_4$ are strong greenhouse gases, and, therefore, must be replaced by the alternative compounds that have low global warming potentials. One possible replacement is $\text{CF}_3\text{I}$ due to its weak C–I bonding. Thus, there are practical as well as fundamental interests in obtaining adequate spectroscopic information of the underlying electronic transitions for these fluorohalomethanes. Though these molecules have been investigated by a variety of techniques, the systematic study is less extensive in low-energy electron spectroscopy (EELS). The present work is concentrated with the electron impact excitation of the threshold electronic states of $\text{CF}_3\text{Cl}$ and $\text{CF}_3\text{Br}$.

In low-energy electron impact spectroscopy, molecular transitions resulting from the transfer of energy from incident electrons to target molecules are not restricted by the optical selection rules. The low-momentum-transfer (high impact energy, low scattering angle) energy-loss spectrum is equivalent to the photoabsorption spectrum. The large-momentum-transfer (low impact energy, high scattering angle) spectra can reveal the presence of optically forbidden excitation. Thus, electron impact spectroscopy is a useful complement to conventional optical spectroscopy in the assignment of molecular electronic states.

In these regards, we will first briefly review the previous works in vacuum ultraviolet (VUV) and electron-impact spectroscopy on primarily $\text{CF}_3\text{Cl}$ and $\text{CF}_3\text{Br}$. Photoabsorption spectra have been studied extensively by Doucet et al.\(^1\) for both $\text{CF}_3\text{Cl}$ and $\text{CF}_3\text{Br}$ in the energy range from 6.2 to 10.33 eV. Gilbert et al.\(^2\) and Doucet et al.\(^3\) obtained the VUV spectra for $\text{CF}_3\text{Cl}$ and $\text{CF}_3\text{Br}$, respectively, in the region extended to 19 eV. Photoelectron spectroscopy using resonance lamps was also carried by Doucet et al.\(^1\) for $\text{CF}_3\text{Br}$ above about 11 eV, and synchrotron radiation measurements by Bozek et al.\(^4\) for $\text{CF}_3X$ ($X=\text{F}, \text{Cl}, \text{Br}, \text{I}$). Biehl et al.\(^5\) reported VUV and visible spectroscopy of $\text{CF}_3X$ ($X=\text{H}, \text{F}, \text{Cl}, \text{Br}$) using fluorescence excitation and dispersed emission technique. Mason et al.\(^6,7\) extensively studied photoabsorption of $\text{CF}_3X$ ($X=\text{Cl}, \text{Br}, \text{I}$) using ASTRID synchrotron source and compared with the results from electron spectroscopy. In the electron-impact studies, zero-angle electron energy loss spectroscopy for $\text{CF}_3\text{Cl}$ was performed by King and McConkey\(^8\) at the impact energy of 500 eV and by Zhang et al.\(^9\) at the impact energy of 8 keV. Momentum-
transfer-resolved electron energy loss spectrum has given information on an unusually low-lying nondipole transition at 7.7 eV in CF3Cl by Ying et al. In addition, their ab initio generalized oscillator strengths (GOS) and single-excitation configuration interaction calculations showed that this transition is consistent with an electronic excitation from the nonbonding lone-pair orbital on the Cl atom \( \text{HOMO} \) to the antibonding \( \pi^* \) or LUMO \( \text{LUMO} \) leading to dissociation of the C–Cl bond. Chen et al. studied the HOMO of CF3Cl by binary \( e, 2e \) electron momentum spectroscopy. They calculated the position density and the result shows that the electron concentrates near the nucleus of the Cl atom implicating that this orbital is the chlorine’s lone-pair orbital. And the momentum density shows that the orbital is \( p \) type. Sunohara et al. measured absolute differential cross sections for elastic and vibrationally inelastic scattering of electrons from CF3Cl and CF3Br using a crossed-beam experiment in the energy range of 1.5–100 eV.

In this report, we have examined systematically the energy-loss region from 3 to 18 eV for all the halofluoromethanes up to CF4 at incident electron energies of 30 and 100 eV and at the scattering angles of 4° or 3° for CF3H and CF3I) and 15° and for incident electron energies of 30 and 100 eV. For comparison purpose, an “optical absorption” type spectrum was taken using high-energy, low-angle conditions. A number of tube lenses in the spectrometer have been used for imaging and energy control of the electron beam, whose characteristics were carefully modeled by electron trajectory calculations. Both the monochromator and analyzer are enclosed in differentially pumped boxes to reduce the effect of background gases and to minimize any stray electron background. Overall resolution in all spectra was 25–35 meV full width at half maximum. The angular scale was calibrated to ±1.5° by noting the symmetry of elastic scattering about the true zero-degree point. The energy scale was calibrated to the He 19.367 eV resonance at 90°.

III. THEORY

The adiabatic potential energy curves for the CF3Cl and CF3Br molecules in their electronic \( X^1\text{A}_1 \) ground state were obtained at the density functional theory level. The three-parameter density functional of Becke was applied with the Lee, Young, and Parr correlation functions (Becke3LYP).
Also, under \( T_d \) symmetry of \( CF_4 \), the following configuration is obtained for the ground state:

\[
(1a_1)^2(1t_2)^6(2a_1)^2(3a_1)^2(2t_2)^6(4a_1)^2(3t_2)^6(1e)^4(4t_2)^6(1t_1)^6
\]

for \( CF_4 \).
besides the energies required for the dissociation are getting lower along the substitution as observed in Fig. 2. Those low-lying states are, for the first time, systematically identified as shown in Fig. 2 for CF$_3$Cl, CF$_3$Br, and CF$_3$I by the low energy electron impact spectroscopy at 30 eV.

The observations are summarized for each fluoromethyl halide as follows.

A. CF$_3$H

Following the previous measurements for the higher impact energies of 400 eV as well as for the photoelectron spectroscopy, the peaks at 10.92 and 11.95 are assigned to $6a_1 \rightarrow 3s$ and $6a_1 \rightarrow 3p$, and also the higher features at 12.58, 13.65, and 14.99 eV to $5e \rightarrow 3s$, $5e \rightarrow 3p$, or $4e \rightarrow 3s$ and $4e \rightarrow 3p$ transitions using term values, respectively. To compare CF$_3$H with CF$_4$, we recall the fluoromethane series CH$_3$F, CH$_2$F$_2$, CHF$_3$, and CF$_4$ including CH$_4$ (Refs. 15 and 16) and discuss the substitution effect of F atom from the CH$_4$ molecule. The highest orbital in CH$_4$ has the $1t_2$ symmetry of a carbon-hydrogen sigma bonding correlated only with the carbon 2$p$ orbitals.$^{17}$ Along the halogen substitution from CH$_4$ to CHF$_3$, however, the uppermost component of the orbital $1t_2$ changes to more or less carbon-fluorine pi-antibonding character. This trend is just that presented by the first band at 10.92 eV in CF$_3$H, but which was observed at lower energy side around the same energies of $\sim$9.42 eV in CH$_4$, CH$_3$F$_2$, and CH$_2$F$_2$. Thus, in CF$_3$H, the first two bands are assignable as $6a_1 \rightarrow 3s$ and $3p$.

B. CF$_4$

The EELS and photoelectron results show that all excitations from the ground state have been classified as Rydberg feature as shown in Fig. 2. The transitions of $1t_1 \rightarrow 3s$ and $1t_1 \rightarrow 3p$ or $5a_1$ correspond to the peaks at 12.6 and 13.7 eV, respectively. The two highest molecular orbitals, $1t_2$ and $4t_2$, are overlapped closely by 1.2 eV.$^{15}$ The former transition is expected to have a symmetry forbidden character due to weak oscillator strength of 0.024±0.003. It is noted that the $1t_1$ orbital is composed of fluorine lone-pair pi orbitals, while $4t_2$ is carbon-fluorine sigma bonding. And, the latter transition indicates more or less carbon-fluorine pi-antibonding character due to HOMO $\rightarrow$ LUMO transition. The next higher features at 13.9 and 15.8 eV are assigned as the transitions of $4t_2 \rightarrow 3s$ and $4t_2 \rightarrow 4s, 3d$ or $1e \rightarrow 3p$, respectively. The band at 14.8 eV is assignable as $1e \rightarrow 3s$, indicating a forbidden transition because this band is clearly observed in more favorable experimental condition at 30 eV for 20° than at 100 eV for 4°.

C. CF$_3$Cl

Figure 3 shows the electron energy loss spectra for CF$_3$Cl at the impact energy of 30 eV and the scattering angle

![Image](https://example.com/image.png)

**FIG. 2.** (Color online) Electron energy loss spectra of CF$_3$X ($X$ = H, F, Cl, Br, I) in the energy-loss region from 3 to 18 eV at incident electron energies of 30 and 100 eV and at scattering angles of 4° (3° for CF$_3$H and CF$_3$I) and 15°, respectively. Solid line represents the data at lower energy/larger angle and dotted line higher energy/smaller angle.

![Image](https://example.com/image.png)

**FIG. 3.** Electron energy loss spectra of CF$_3$Cl in the energy-loss region from 4 to 14 eV at incident electron energies of 30 and 100 eV and at scattering angles of 15° and 4°, respectively.
of 15° and at the impact energy of 100 eV and the scattering angle of 4°. Based on numerous results from the photoelectron and the dipole EELS measurements, in the valence-shell electronic structures of CF₃Cl, the peaks at 9.6, 10.6, and 11.5 eV are assigned to 6e → 4s, 6e → 4p, 6e → 5s, 3d, respectively, and peaks 11.5 and 13.2 eV are also identified as 7a₁ → 4s, 7a₁ → 3d, 4d, 5s, respectively, terminating to the Rydberg limit of 7a₁⁻¹. 6e (HOMO) orbital is comprised essentially of a lone pair on Cl atom, and 7a₁ is an orbital associated with C–Cl bonding. Therefore, the first ionization potential near 13.0 eV corresponds to taking out of one electron from the chlorine lone pair orbital, and the second ionization potential near 15.2 eV is related to the orbital populated in the C–Cl bond.2 As shown in Fig. 3, around 7.7 eV, the peak with a low-lying nondipole character is clearly emerged at the impact energy of 30 eV for the scattering angle of 15° but not at the impact energy of 100 eV for 5°. This nondipole electronic transition with a quadrupole characteristic is observed for the first time in the electron energy loss experiments. However, it was previously observed by the different experimental method using the momentum-transfer-resolved electron loss spectroscopy10 at 2.5 keV, which was assigned as 6e(HOMO) → 8a₁(LUMO) transition related an electronic transition from a nonbonding Cl 3p orbital (HOMO) to an antibonding 5σ₃− orbital (LUMO). In the case of n(C–Cl) → σ*:C–Cl transition, the overlap between the initial-state and final-state orbital wave functions occurs mostly at Cl site. The transition moment of this electronic transition is expected to resemble that of p-to-p transition in Cl atom which explains the quadrupole nature.10,18–20 In the study by Ying et al.,10 the GOS of this transition has been determined and found to have a shape characteristic of a quadrupole transition with a maximum at around 1 a.u. of the momentum transfer. While similar low-lying quadrupole transitions have been observed for other members of CFC, no such nondipole transition is found in the case of CF₃, suggesting the involvement of Cl atoms in this quadrupole feature. The LUMO was found to be a repulsive potential character by McLean and Chandler.21

D. CF₃Br

The electron energy loss spectrum for CF₃Br is given in Fig. 4. As assigned previously from the UV spectrum data, the strong bands at 8.7 and 9.53 eV are associated with the Br 6e⁻¹ ionization at 12.0 eV, the transitions from the bromine lone pair orbital to 5s and 5p terminating orbitals, respectively. The former one has a shoulder at 8.42 eV, which is due to spin-orbit splitting and there is no corresponding shoulder observed for CF₃Cl, while this splitting is much clearer for CF₃I than for CF₃Br. The peak at 8.7 eV is a singlet-singlet transition and the shoulder at 8.42 eV has a singlet-triplet character.1 However, Eden et al.7 argue that the relative intensity of the two peaks at 8.42 and 8.7 eV are approximately equal in 100 eV, 4° and 30 eV, 15° spectra (see Fig. 2) and this lack of angular dependence strongly suggests that both features are due to singlet to singlet transitions. This leads Eden et al. to suggest that the broad feature peaked at 8.42 eV is due to the promotion of an electron from the second highest occupied molecular orbital (MO) (C–Br σ) to the LUMO (C–Br σ*). The band at 10.54 eV corresponds to the transition of 7a₁ → 5s. The first ionization potential corresponds to Br lone pair orbital, and the second ionization potential belongs to bonding orbitals of mainly C–Br character. Compared to CF₃Cl, the ionization potentials of 6e and 7a₁ for CF₃Br are lower by around 1 eV.5 Though, as mentioned above, no direct comparison is possible in the intensity, the minimum and the peak positions are apparently changed in the spectral shapes around 9.53 eV due to the optical forbidden transitions involved strongly. The most prominent feature is a clear observation again for the low-lying electronic excitation states near threshold around 6.5 eV at electron impact energy only of 30 eV not of 100 eV. But, this feature is not mainly due to the optical forbidden character, like a triplet state, because the weak transition has been also observed in the vacuum ultraviolet spectrum.1

E. CF₃I

The assignments for the transitions are reported previously (Ref. 6 and references therein). Briefly, the weak broad band observed around 4.6 eV was identified as 5pπ → σ* (C–I) transition which is composed in the 3Q₁ (4.10 eV), 3Q₀ (4.70 eV), and 1Q₁ (5.34 eV) states due to the spin-orbit coupling and exchange interactions as confirmed by the photoabsorption result. The 3Q₀ and 1Q₁ states are enhanced at 4.70 and at 5.34 eV for the impact energy of 30 and 100 eV, respectively, as shown in Fig. 2. Moreover, the resultant excited states, i.e., the broad features without a fine vibrational progression, are leading to dissociation along the C–I stretch coordinate due to occupation of the C–I (5’S₁/₂, 5’S₇/₂) antibonding state, respectively, where 5pπ is the outermost lone-pair π orbital of I atom and σ* is antibonding carbon iodide sigma MO. The peaks at 7.13 and at 7.75 eV are assigned as
6e_{3/2} \rightarrow 6s \text{ and } 6e_{1/2} \rightarrow 6s \text{ Rydberg transitions, giving rise to four components: } 2^2E_{3/2}[2], 2^2E_{3/2}[2], 2^2E_{1/2}[0], \text{ and } 2^2E_{1/2}[1] \text{ again due to the spin-orbit and exchange interactions. On those bonds with the bound state nature, the weak vibrational progressions are visible even with the present energy resolution of our spectrometer though the bands exhibit clearly structure dominated by the vibrational progressions in their high resolution photoabsorption spectra. The further bands being observed at 8.8 and 9.6 eV are tentatively attributed to the transition of } 6e_{3/2} \rightarrow 7s \text{ and } 6e_{1/2} \rightarrow 7s, \text{ and weak features observed between 8.5 and 8.6 eV may be due to a vibrational progression and the band between 9.1 and 9.4 eV due to excitation of yet another state. There is no information available for a series of the bands above 10 eV, and thus a full electronic structure calculation deriving the energies of high-lying Rydberg series would be desired.}

V. CONCLUSIONS

By varying the scattering angle and incident energy, we have studied systematically a low-lying transition in the valence shells of the halogen-substituted methane molecules CF_{3}X (X=H,F,Cl,Br,I). As mentioned above, those low-lying states are, for the first time, systematically identified for CF_{3}Cl, CF_{3}Br, and CF_{3}I by the low energy electron impact spectroscopy at 30 eV, in which the large-momentum-transfer (low impact energy, high scattering angle) spectra are, in general, favorable to detect the presence of optically forbidden excitations. Furthermore, including CF_{4a}, these all four molecules have a lone-pair electron in each halogen atom, but not in the halogen atom of CF_{3}H. From the molecular-symmetry point of view, former three belong to C\text{\_3p}, but CF_{4a} belongs to a different and much higher symmetry of T\text{\_d}. Therefore, it is concluded that the halogenated CF_{3}Cl, CF_{3}Br, and CF_{3}I with C\text{\_3p} symmetry as well as a lone-pair electron reveal the low-lying weak band, leading to dissociation of the molecules into the CF_{3} radical along the C–X molecular axis. Furthermore, the CIS molecular calculation reveals the repulsive potential nature of LUMO for those two molecules, which supports the present experimental results. The direct product table for the symmetry group C\text{\_3p} indicates that the 1a_{1} \rightarrow 1e \text{ transition is both dipole allowed and quadrupole allowed, and the small transition dipole moment therefore suggests that this transition can be attributed predominantly to quadrupole interactions. In addition, as indicated by the dominant molecular orbital coefficients, the 1a_{1} \rightarrow 1e \text{ transition corresponds to a transition from an essentially nonbonding Cl 3p orbital to an antibonding p–\sigma^{*} \text{ C–Cl orbital. The p-to-p transition is therefore consistent with a quadrupole transition. It should also be noted that a vertical transition originated from the ground vibrational level of the } ^{1}A_{1} \text{ state terminates at the repulsive part of the potential energy curve of the } ^{1}E \text{ state above its dissociation limit.}

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