A REVISED PLANETARY NEBULA LUMINOSITY FUNCTION DISTANCE TO NGC 628 USING MUSE

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ABSTRACT

Distance uncertainties plague our understanding of the physical scales relevant to the physics of star formation in extragalactic studies. The planetary nebulae luminosity function (PNLF) is one of very few techniques that can provide distance estimates to within $\sim 10\%$; however, it requires a planetary nebula (PN) sample that is uncontaminated by other ionizing sources. We employ optical integral field unit spectroscopy using the Multi-Unit Spectroscopic Explorer on the Very Large Telescope to measure [O III] line fluxes for sources unresolved on 50 pc scales within the central star-forming galaxy disk of NGC 628. We use diagnostic line ratios to identify 62 PNe, 30 supernova remnants, and 87 H II regions within our fields. Using the 36 brightest PNe, we determine a new PNLF distance modulus of $29.91^{+0.08}_{-0.13}$ mag ($9.59^{+0.35}_{-0.57}$ Mpc), which is in good agreement with literature values, but significantly larger than the previously reported PNLF distance. We are able to explain the discrepancy and recover the previous result when we reintroduce SNR contaminants to our sample. This demonstrates the power of full spectral information over narrowband imaging in isolating PNe. Given our limited spatial coverage within the Galaxy, we show that this technique can be used to refine distance estimates, even when IFU observations cover only a fraction of a galaxy disk.

Key words: galaxies: distances and redshifts – galaxies: individual (NGC 628) – ISM: supernova remnants – planetary nebulae: general

1. INTRODUCTION

When connecting star formation and galaxy evolution, nearby galaxies ($D < 20$ Mpc) benefit from an external view that is still accessible at high physical resolution (10–100 pc), using ground based instruments ($\sim 1''$ seeing). However, a major source of uncertainty in the study of nearby galaxies is their distances. Reliable methods for distance determination ($< 5\%$ accuracy), based on Cepheids and the tip of the red giant branch (TRGB) measurements, are challenging to observe at large ($D > 5$ Mpc) distances, requiring the resolution of individual stars (Jacoby et al. 1992). SNe II (Dessart & Hillier 2005; Kasen & Woosley 2009) and Tully–Fisher (Freedman & Madore 2010) methods, often used in nearby galaxies, suffer from much larger ($\sim 20\%$) uncertainties. The [O III] $\lambda 5007$ Å planetary nebula luminosity function (PNLF) provides an alternate, yet still accessible, method for obtaining accurate ($< 10\%$) distances (Ciardullo et al. 1989; Jacoby et al. 1992; Ciardullo 2010, 2013).

Planetary nebulae (PNe) are ionized by low and intermediate mass stars ($1\text{--}8 M_{\odot}$), with ages 0.1–1 Gyr, as they evolve from the asymptotic giant branch to the white dwarf phase (Paczynski 1971; Iben & Renzini 1983). The central stars are very bright ($> 6000 L_{\odot}$; Vassiliadis & Wood 1994), with $\sim 12\%$ of the total luminosity of the central star reprocessed into the [O III] $\lambda 5007$ Å emission line (Dopita et al. 1992; Schönberner et al. 2007, 2010), making them relatively easy to observe in emission line surveys. The PNLF exhibits a sharp exponential cutoff, making the top ~ 1 mag of the luminosity function well-

suited for use as a standard candle (Ciardullo et al. 1989). Although theoretical models predict dependences on the stellar population age and metallicity (Dopita et al. 1992; Marigo et al. 2004), the PNLF has been empirically shown to be invariant across a variety of galaxy types (elliptical, spiral and irregular) that host very different underlying stellar populations (Ciardullo et al. 1989; Feldmeier et al. 1996, 1997).

Previous PNLF studies employed narrowband imaging, which is strongly contaminated by background stellar continuum emission from the Galaxy. As a result, most of the identified PNe in these surveys are located at large radii (Herrmann et al. 2008), where the contamination is reduced, but the PN density (which traces the stellar surface density) is lower (Ciardullo et al. 1989). Optical integral field spectroscopy provides full spectral information, allowing us to cleanly fit and remove emission from the stellar continuum. In star-forming spiral galaxies, this technique also allows us to model and remove background line emission arising from the diffuse interstellar medium (Kreckel et al. 2016), and provides simultaneous observation of a suite of bright emission lines (H α , [N II] $\lambda 6584$ Å, [S II] $\lambda 6717$ Å, and [S II] $\lambda 6737$ Å) that can be used to distinguish PNe from other [O III] emitters, such as H II regions and supernova remnants (Sabbadin et al. 1977; Baldwin et al. 1981; Riesgo & López 2006). Past integral field unit (IFU) studies have demonstrated the feasibility of this IFU technique when extracting PN line fluxes (Roth et al. 2004; Sarzi et al. 2011), but lacked a large field of view and observed only a handful of PNe. Recent and ongoing IFU surveys of nearby galaxies are typically designed to trace star formation

and measure stellar populations, but in many cases will have sufficient depth and spatial resolution to simultaneously detect a large population of PNe, providing significant improvements on the distance estimates for these galaxies.

We examine two regions within the nearby grand design spiral galaxy NGC 628, using Very Large Telescope (VLT)/Multi-Unit Spectroscopic Explorer (MUSE) optical IFU spectroscopy that was obtained in order to study the contrast between spiral arm and interarm star formation (Kreckel et al. 2016). Its face-on orientation ($i = 9^\circ$; Blanc et al. 2013b) minimizes cross-contamination of PNe with other objects, and reduces the impact of internal dust extinction on our observations. NGC 628 has a well-studied metallicity gradient of $12 + \log(\text{O}/\text{H}) = 8.834 - 0.485 \times R$ (dex R_{25}^{-1}), as measured using direct abundances based on observations of the temperature-sensitive auroral lines (Berg et al. 2015), with roughly solar abundances (Asplund et al. 2009) within the central ($R < R_{25}$) region. This galaxy has an uncertain distance, with estimates ranging from 7 to 10 Mpc—based on a variety of techniques (see Section 4.1), including a narrowband PNLF study (Herrmann et al. 2008).

In Section 2, we present our MUSE data. In Section 3, we identify PN candidates and, using diagnostic line ratios, exclude H II region and SNR contaminants from the sample. In Section 4, we calculate a new PNLF distance to NGC 628, compare it to previous distance estimates, explain our discrepancy with the previous PNLF distance, and explore what the limiting distance is for our technique. We conclude in Section 5.

2. DATA

We observed NGC 628, using the MUSE spectrograph (Bacon et al. 2010) at the VLT. This powerful new optical IFU provides a $1' \times 1'$ field of view with $0''.2$ pixels and a typical spectral resolution of $\sim 2.75 \text{ \AA}$ (150 km s^{-1}). We observed NGC 628 in three positions: two northern (ID 094.C-0623) and one southern (ID 095.C-0473), with all data reduction details provided in Kreckel et al. (2016). Total on-source exposure times are 42 (50) minutes for our northern (southern) pointings. Observations were taken using the nominal wavelength range, covering 4800–9300 Å. Typical seeing is $0''.8$ across all fields. Given the range of distances reported for this galaxy, this corresponds to a spatial resolution of 30–50 pc.

We validate the astrometry and flux calibration of the resulting MUSE data cubes against SDSS r-band images (York et al. 2000). We construct a simulated r-band image from our data cube by convolving the r-band filter shape with the MUSE spectra at each spatial position. We then fit compact sources in both the simulated MUSE r-band image and the SDSS r-band image, using SExtractor (Bertin & Arnouts 1996). We find the source positions agree within $0''.2$ across all fields. To check the absolute flux calibration, we sample both images with randomly placed $5''$ apertures, and find systematic agreement (Figure 1) to within 1% and a scatter of less than 5% (0.05 mag), confirming the high level of accuracy in our MUSE data set.

Following Kreckel et al. (2016), we fit the stellar continuum with MIUSCAT templates (Vazdekis et al. 2012) and measure emission lines using LZIFU (Ho et al. 2016) assuming single Gaussian fits. These emission lines can serve as diagnostics for the mechanism that ionizes and excites the gas. Figure 2 presents a three-color image of all fields, combining the [O III],

Figure 1. Comparison of MUSE with SDSS photometry in the r-band for $5''$ apertures across all fields. Fluxes systematically agree to within 1% (solid line shows a one-to-one agreement) with a scatter of less than 5% (0.05 mag). Outliers in the southern field lie at the edge of bright star clusters where we are sensitive to slight ($<0''.2$) offsets in astrometry.

$\text{H}\alpha$, and [S II] emission lines, highlighting the difference in color between the star-forming complexes and surrounding diffuse ionized gas, as well as revealing many unresolved objects bright in [O III], which are likely PNe (see Section 3.1).

Because the PNLF is quite sensitive to the absolute flux calibration, we further check our measured line fluxes against the line fluxes observed by the optical IFU survey VENGA (Blanc et al. 2013a, 2013b). Here, we have convolved our MUSE line maps to match the $5''.5$ resolution in VENGA and compare the [O III] and $\text{H}\alpha$ line fluxes (Figure 3). We find our line fluxes agree within 3%, with a scatter of 20%. For this comparison, we exclude regions that show the largest discrepancy, because they are at the position of an overlapping bright H II region and a bright foreground star, which fall at the edge of the field of view in both images. The combination of edge effects and significant point-spread function wings could bias our convolution, thus invalidating our comparison near this region. The overall agreement between the emission lines maps observed with these two very different instruments confirms that the systematics within the MUSE and VENGA data cubes are minor, and demonstrates the overall high quality of IFU observations possible in nearby galaxies.

3. RESULTS

We identify our PN candidates based on unresolved objects in the [O III] emission line maps. However, we must also clean our candidate list of contamination by other line emitting objects, such as H II regions and supernovae remnants, before fitting the PNLF to determine a distance to the Galaxy.

3.1. Identifying PN Candidates

Because bright PNe in the Galaxy are much smaller than our 30–50 pc ($0''.8$) spatial resolution (<1 pc diameter; Acker

Figure 2. V-band image of NGC 628, overlaid with three color images for the MUSE fields, using three emission line maps: [O III] (blue), H α (green) and [S II] (red). These emission lines serve as diagnostics for the physical conditions in the ionized gas. PN candidates appear as compact blue objects.

et al. 1992), we use the IDL task FIND (an adaptation of DAOPHOT) to identify unresolved objects within the [O III] line maps (Figure 4). We perform aperture photometry at each position for apertures $1''.6$ (two times the FWHM) in diameter. As our field of view is entirely contained within the disk of the Galaxy, we do not have any stars suitable for a standard curve of growth analysis to correct for lost flux outside of the aperture and robustly determine the seeing. For this reason, a slightly sub-optimal aperture size was chosen to conservatively ensure most of the flux is included in the aperture. In addition, because our line fluxes are measured from spectral fitting—which does not allow for negative fluxes—our line maps are biased to positive fluxes in regions of low signal-to-noise. This positive bias is not a background, and impacts only regions close to our detection limits. Our apertures are small enough that this is not a concern when totaling the flux within the aperture; however, it complicates any modeling of the aperture bias when attempting to recover flux from the wings of the PSF.

There are no foreground stars within our field suitable for performing the aperture correction. Therefore, we perform our curve of growth analysis on a model PSF, using parameters obtained from the brightest point source [O III] line emitters. This also avoids the bias to positive fluxes at regions with low

signal-to-noise. We convert these aperture corrected fluxes, $F_{[\text{O III}]}$, to an apparent magnitude as

$$m_{[\text{O III}]} = -2.5 \log F_{[\text{O III}]} - 13.74 \quad (1)$$

where $F_{[\text{O III}]}$ is given in $\text{ergs cm}^{-2} \text{s}^{-1}$ (Jacoby 1989).

All objects are extinction-corrected for the Milky Way foreground emission assuming $E(B - V) = 0.062$ (Schlafly & Finkbeiner 2011), $R_V = 3.1$ and the Cardelli et al. (1989) extinction law ($A_V \sim 0.2$ mag). We apply no internal extinction correction based on the results of Feldmeier et al. (1997), who argue that, for face-on galaxies, the expected extinction in $m_{[\text{O III}]}$ for face-on galaxies is less than 0.1 mag. This is supported by Herrmann et al. (2008), who found that the PNLFs for the inner and outer disk of NGC 628 are the same. As dust profiles are typically exponentially decreasing, with little extinction suffered in the outer disk (Giovanelli et al. 1994; Muñoz-Mateos et al. 2009), the lack of radial variation in the PNLF suggests that, for NGC 628, there is also minimal extinction in central regions. Finally, we observe reddening consistent with only the Milky Way foreground ($A_V = 0.21 \pm 0.15$ mag) from our stellar spectral fitting, further supporting our decision not to apply any corrections for internal extinction.

Figure 3. Comparison of the $H\alpha$ and $[O\text{ III}]$ line maps obtained in this study to those from VEGA (Blanc et al. 2013b). Line fluxes agree within 3% (solid line shows a one-to-one agreement), with 20% scatter (corresponding to 0.2 mag in the $[O\text{ III}]$ line). The large, order-of-magnitude offset observed in both lines toward high fluxes is at the position of an overlapping bright foreground star and H II region (diamonds), which also falls toward the edge of both the MUSE and VEGA field of view. This combination of edge effects and extreme PSF wings invalidates our comparison near this region.

Herrmann et al. (2008) surveyed NGC 628 using $[O\text{ III}]$ narrowband imaging and identified 153 PNe. 12 PNe are located within our field of view, and for these, we compare $m_{[O\text{ III}]}$ with their reported values (Figure 5). We find reasonably good agreement between the two samples; most agree within 0.2 mag, close to our typical uncertainty. We find astrometric agreement within $1''$ in all cases, with typical offsets of $0''.5$.

Based on our estimated uncertainties, we are confident in our detections down to $3\sigma = 27.0$ mag in the northern fields and, due to the longer exposure time and slightly better observing conditions, 27.5 mag in the southern field. Due to the relatively low and uniform background (Figure 4), we expect we are complete and robustly detect all sources down to these limits.

3.2. Removing Contaminants

Our $[O\text{ III}]$ detected sample contains not just PNe, but also compact H II regions or recent supernovae remnants (SNRs). In order to identify contaminants to our PN sample for all $[O\text{ III}]$ detections, we also measure line fluxes within the same aperture from $H\alpha$, $[N\text{ II}] \lambda 6584 \text{ \AA}$, $[S\text{ II}] \lambda 6717 \text{ \AA}$, and $[S\text{ II}] \lambda 6737 \text{ \AA}$. Combination of these emission lines provides diagnostic ratios that allow us to identify the most likely ionizing source for each of our $[O\text{ III}]$ detections. Background galaxies, which contaminate narrowband imaging surveys, are ruled out here by our simultaneous detection of multiple emission lines for all objects.

More exotic ionizing sources are also possible. Symbiotic binaries would be detectable through a very red stellar continuum spectrum, which we do not see for any of our sources. We find three objects (Table 1) with very broad Balmer emission lines ($>1000 \text{ km s}^{-1}$), strong He I emission, and the Ca II triplet in emission. Balmer lines in some sources

show P Cygni profiles, and all sources fall in or near large star-forming complexes. These features are all consistent with Wolf-Rayet stars (Crowther 2007), and we exclude these objects from our analysis.

Our aperture photometry may be biased by the background line emission arising from the diffuse ionized medium that is unrelated to the compact $[O\text{ III}]$ sources (Madsen et al. 2006; Haffner et al. 2009; Kreckel et al. 2016). Detailed two-dimensional modeling of the diffuse background is necessary to ensure its careful removal from the line fluxes associated with the compact source (Roth et al. 2004). However, because we are mainly interested in using these lines to identify contaminants in our PN sample, we employ a simple median smoothing over $2''$ scales (roughly three times the FWHM) to identify and remove the local diffuse background emission at the position of each $[O\text{ III}]$ detection. Masking discrete sources does not significantly change our results. This background subtraction is only applied in cases where the median smoothed background is detected with a signal-to-noise ratio greater than 5, which applies to most of the $[N\text{ II}]$, $H\alpha$, and $[S\text{ II}]$ line maps, but only 4% of the $[O\text{ III}]$ maps. We identify three PN candidates in positions where there is a high background in $[O\text{ III}]$ and the detection lies superimposed on a star-forming region. Due to the difficulties in accurately subtracting the background, and the increased likelihood of these being contaminants, due to their position, we omit these three objects from our analysis. As the PNe are expected to be uniformly distributed throughout the disk, exclusion of this 4% of our field of view should not bias our luminosity function.

3.2.1. H II Regions

For narrowband PN searches, $H\alpha$ imaging is commonly used to remove H II regions, which are the main contaminant to our

Figure 4. [O III] emission line maps for our MUSE fields. Identified PNe (white circles) are uniformly distributed throughout the fields, consistent with their association with a moderately old (10^8 – 10^9 year) stellar population. Compact SNRs (red circles) are a significant potential contaminant to our PN sample.

Figure 5. A comparison of measured [O III] apparent magnitude for 12 overlapping PNe observed in a previous narrowband imaging study (Herrmann et al. 2008; Herrmann & Ciardullo 2009). We find very good ($\sim 1\%$) systematic agreement between the two samples (solid line represents a one-to-one agreement), with most agreeing within 0.2 mag (20%), close to our typical uncertainty.

Table 1
Emission Line Stars

Name	R.A. (J2000)	Decl. (J2000)
1	01:36:47.85	+15:47:10.1
2	01:36:44.17	+15:47:16.0
3	01:36:45.16	+15:45:02.9

sample. At the gas-phase metallicity of NGC 628, H II regions have $H\alpha$ flux brighter than [O III] (Shaver et al. 1983), whereas the inverse is true in PNe (Baldwin et al. 1981). Ciardullo et al.

(2002) showed that PNe typically have line ratios of

$$4 > \log \frac{[\text{O III}]}{H\alpha + [\text{N II}]} > -0.37M_{[\text{O III}]} - 1.16, \quad (2)$$

where this has been empirically determined using the $H\alpha$ + [N II] line due to contamination of [N II] within the narrowband filters. Other diagnostic diagrams are available, based on spectroscopic observations rather than narrowband imaging, to distinguish different ionizing sources (Sabbadin et al. 1977; Baldwin et al. 1981; Riesgo & López 2006). However, due to the natural variations between PNe physical conditions, and given that we have only [O III] and $H\alpha$ detected for many of our PNe, we find this empirical narrowband criteria the most effective. Assuming a 50% contribution by [N II] (Riesgo & López 2006) and a distance modulus of 29.91 mag (see Section 3.3), we apply Equation (2) to our sample to identify 87 H II region contaminants (Figure 6). Assuming a lower [N II] contribution (Sabbadin & D’Odorico 1976; Baldwin et al. 1981) or varying the distance modulus within the uncertainty range does not significantly affect the resulting PNLF and distance estimate. Many of our sources are not detected in $H\alpha$, and thus have only lower limits. In addition, many sources appear to achieve ratios well above 4; however, due to the large uncertainties in the $H\alpha$ line fluxes, we retain these sources as potential PN candidates.

3.2.2. Compact Supernova Remnants

Another contaminant that is more difficult to identify and remove is compact supernova remnants (SNRs). While the most evolved SNRs can span ~ 100 pc, and are thus expected to be resolved by our survey, younger remnants are only ~ 20 pc in size and will remain unresolved (Franchetti et al. 2012). None of the three SN known within NGC 628 (SN 2002ap, SN 2003gc, SN 2013ej) are within our observed fields, however previous study of diagnostic line ratios in SNR, H II region and PN populations (Sabbadin et al. 1977; Riesgo & López 2006)

Figure 6. $[\text{O III}]/\text{H}\alpha$ line ratios as a function of $[\text{O III}]$ apparent magnitude for our full sample of $[\text{O III}]$ detected objects, including lower limits where $\text{H}\alpha$ is not detected at a 3σ level. We also show 1σ uncertainties. PNe (filled circles) are distinguished from H II regions (plus symbols) following the diagnostic (dotted line) determined by Ciardullo et al. (2002) using local group PNe. SNRs (open red circles) identified using additional line ratio diagnostics would significantly contaminate our PNe sample using only this threshold. All marked PNe form our final sample of PNe (Table 3).

reveal that they inhabit very different regions of a diagram that compares $\text{H}\alpha/[\text{N II}]$ and $\text{H}\alpha/[\text{S II}]$ line ratios (Figure 7). Many of our PN candidates are not detected in any of these lines, however we can still use the line ratio when detected to identify contaminants to our sample.

We find that the H II regions identified in Figure 6 are consistent with the expected region of the diagram, and find a population of sources with line ratios consistent with SNRs. Due to the difficulties in accurately subtracting the diffuse background emission from the emission lines other than $[\text{O III}]$, and the faintness of these lines, the uncertainties on our line ratio measurements are large; therefore, we employ a simple threshold to identify all sources with detected line ratios $\text{H}\alpha/[\text{S II}] > 2.5$ as SNRs (Blair & Long 2004). A few of our PN candidates have line ratio detections or limits that place them in regions typically populated by H II regions or SNRs. However, given that these sources are not inconsistent with PNe (within the uncertainties), we retain them in our sample.

We identify a total of 30 SNR candidates (Figure 4; Table 2), including one object (#17) previously identified as a PN (M74-30) by Herrmann & Ciardullo (2009). Our fields do not cover the six SNRs previously identified from narrowband imaging (Sonbaş et al. 2010). We do not expect this sample to necessarily be complete because we are identifying only those SNRs that are compact enough to be confused with PNe. Because the SNRs with the highest $[\text{O III}]$ are expected to have strong shocks, we measure the velocity dispersion in the $[\text{O III}]$ line and find a median value of $\sim 200 \text{ km s}^{-1}$, slightly resolved given the instrumental dispersion of 150 km s^{-1} and further confirming our identification of these sources as SNRs. This

Figure 7. $\text{H}\alpha/[\text{N II}]$ as a function of $\text{H}\alpha/[\text{S II}]$ for all sources with 3σ detections in these emission lines. 1σ uncertainties are shown. The zones associated with H II regions, SNRs, and PNe proposed by Sabbadin et al. (1977) are marked, however more recent work shows that PNe populate a broader range than this simple diagnostic line would suggest (Riesgo & López 2006). As described in the text, we identify SNRs (open red circles) and H II regions (plus symbols) by their position in the diagram. No PNe (filled circles) have detections in all three emission lines, and although those shown have lower limits on $\text{H}\alpha/[\text{S II}]$ consistent with SNRs or H II regions, we retain them in our PN sample because they are not inconsistent with PNe. All marked PNe are in our final sample of PNe (Table 3).

includes the SNR #17, previously identified as a PN, for which we measure a velocity dispersion of $\sim 250 \text{ km s}^{-1}$. The median velocity dispersion for our PNe sample, with expected intrinsic line widths of $\sim 20 \text{ km s}^{-1}$ (Richer et al. 2010a, 2010b), is consistent with being unresolved. Also due to the shock excitation, SNRs are expected to have strong $[\text{O I}]$ compared to $[\text{O III}]$, in contrast to PNe where $[\text{O I}]$ should be 100 times fainter (Baldwin et al. 1981). Our data is not deep enough to place strong limits on the $[\text{O I}]$ flux; however, we detect $[\text{O I}]$ for 17 of our SNR, but only one PN (#13, which may indicate it has been misclassified). These SNRs significantly contaminate our sample, as determined by our $[\text{O III}]/\text{H}\alpha$ diagnostic (Figure 6). This has significant implications for the derived PN luminosity function and measured distance (see Section 4.1).

3.3. Calculating the Distance

After applying our line diagnostic tests, we identify 63 PNe within our MUSE fields (Figure 4; Table 3), including the 11 PNe previously identified by Herrmann & Ciardullo (2009) that fall within our field of view. We reclassify one additional PN from their catalog as an SNR (see Section 3.2.2).

The PNLF has been empirically shown to be well-fit by an equation of the form

$$N(M) \propto e^{0.307M} (1 - e^{3(M^* - M)}), \quad (3)$$

where M is the absolute magnitude and M^* is the absolute magnitude of the most luminous PN (Ciardullo et al. 1989;

Table 2
SNR Identifications

Name	R.A. (J2000)	Decl. (J2000)	$m_{[\text{O III}]}$ (mag)	[O III]/H α	H α /[S II]	H α /[N II]
1	01:36:45.14	+15:45:23.3	25.29	0.46	1.40	1.60
2	01:36:47.65	+15:47:06.5	25.30	0.88	2.17	2.62
3	01:36:41.98	+15:47:11.9	25.44	0.62	0.88	0.72
4	01:36:43.65	+15:47:20.2	25.47	0.85	0.69	0.56
5	01:36:46.24	+15:47:39.2	25.49	2.36	2.32	2.20
6	01:36:44.33	+15:45:27.2	25.54	0.76	2.01	1.88
7	01:36:43.35	+15:47:17.4	25.77	8.94	0.73	0.85
8	01:36:44.07	+15:44:58.2	25.82	>15.82	<0.16	...
9	01:36:42.09	+15:44:56.0	25.83	1.59	2.13	3.15
10	01:36:43.71	+15:47:36.2	25.83	1.36	0.66	0.58
11	01:36:47.40	+15:47:36.0	25.85	2.09	1.04	1.38
12	01:36:42.74	+15:44:43.7	25.97	2.79	1.33	2.01
13	01:36:47.22	+15:47:10.2	25.98	>9.57	<0.16	...
14	01:36:45.81	+15:47:42.1	26.03	3.09	0.40	0.45
15	01:36:45.70	+15:47:39.5	26.07	3.16	0.50	0.63
16	01:36:42.55	+15:45:33.9	26.10	1.36	0.85	0.75
17 ^a	01:36:41.76	+15:47:57.4	26.15	2.47	0.70	0.68
18	01:36:43.89	+15:47:57.4	26.19	1.17	0.50	0.44
19	01:36:43.94	+15:47:15.1	26.26	4.03	0.28	0.23
20	01:36:42.75	+15:45:05.5	26.35	0.59	0.84	1.24
21	01:36:45.30	+15:45:05.1	26.36	1.36	0.96	1.73
22	01:36:42.77	+15:48:18.1	26.51	1.54	1.59	2.42
23	01:36:47.17	+15:47:36.5	26.55	0.73	1.94	1.25
24	01:36:42.91	+15:45:54.5	26.57	3.92	0.81	0.88
25	01:36:41.54	+15:48:04.8	26.63	1.30	0.65	0.68
26	01:36:42.25	+15:47:50.5	26.85	2.48	1.07	0.91
27	01:36:44.86	+15:47:36.7	27.06	1.00	0.53	0.69
28	01:36:43.08	+15:47:56.0	27.09	0.57	1.18	1.29
29	01:36:41.35	+15:48:03.7	27.20	>3.12	<0.63	...
30	01:36:44.20	+15:47:55.3	27.27	0.69	1.37	1.51

Note.

^a Identified as a PN in Herrmann & Ciardullo (2009).

Jacoby et al. 1992). We adopt here the value of $M^* = -4.47^{+0.02}_{-0.03}$ mag, which applies for solar metallicity galaxies such as NGC 628 (Ciardullo et al. 2002). Revised calibrations suggest a slightly higher value ($M^* = -4.53 \pm 0.06$; Ciardullo 2012); however, this would change our distance by only $\sim 3\%$ and is consistent within the uncertainties. The PNLFF cutoff is observed to be fainter in low-metallicity systems, consistent with theoretical models (Dopita et al. 1992; Schönberner et al. 2010). Using the method of maximum likelihood (Ciardullo et al. 1989), we determine the distance modulus that best fits the PNLFF to our observations. This method accounts for the decreased probability of observing PNe near the luminosity cutoff M^* , which is especially important given our small sample size, and removes any biases introduced by binning. We include in our fit only the 36 PNe brighter than 27 mag (Figure 8), our expected completeness limit, and measure a distance modulus of $29.91^{+0.08}_{-0.13}$ mag ($9.59^{+0.35}_{-0.57}$ Mpc). Here, our uncertainty includes (added in quadrature) the uncertainty in the fit ($29.91^{+0.06}_{-0.12}$ mag), our absolute photometric uncertainty (~ 0.03 mag), and a Monte Carlo sampling of our measured apparent magnitude (~ 0.05 mag). Our low number statistics make it difficult to judge the quality of the PNLFF fit; however, the good agreement between our fit and observed PNe sample is apparent when examining the cumulative PNLFF (Figure 9; Mendez et al. 1993).

4. DISCUSSION

Our revised PNLFF distance is significantly larger than that previously determined using narrowband imaging, but consistent with distances obtained through other methods. We test how the careful exclusion of SNRs affects our distance estimate, and show that this contamination can significantly bias our result. Because IFU observations are crucial to isolating a clean PN sample, we further explore the potential for using PNe as distance estimators in ongoing and future IFU studies.

4.1. Comparison with Previous Work

A previous narrowband imaging survey of PNe in NGC 628 measured a distance modulus of $29.67^{+0.06}_{-0.07}$ mag (Herrmann et al. 2008), significantly closer than our measurement. In our comparison with their catalog (see Section 3.1), we identified 12 of their PNe that fall within our field of view and find no systematic offset between the measured [O III] apparent magnitudes (Figure 5). However, following our source classification using diagnostic line ratios (Section 3.2), we reclassify one of their PNe as a SNR. Given that, in Figure 6, we find that 26 SNRs would be classified as PNe without the additional emission line diagnostics provided by our IFU data, we test how the introduction of these SNRs into our PNLFF biases our resulting distance estimate. If we fail to clean the SNR contaminants from our sample, our PNLFF fit is still very

Table 3
Planetary Nebula Identifications

Name	R.A. (J2000)	Decl. (J2000)	$m_{\text{[O III]}}$ (mag)	[O III]/H α	H α /[S II]	H α /[N II]
1	01:36:43.32	+15:47:18.1	25.55	>14.29	...	...
2 ^a	01:36:43.37	+15:47:18.3	25.58	12.67	>0.77	>1.35
3	01:36:43.13	+15:45:54.5	25.69	1.46	3.16	2.53
4	01:36:46.16	+15:47:39.5	25.73	3.53	>2.41	1.73
5 ^a	01:36:44.11	+15:47:20.2	26.10	4.64	>1.31	>2.29
6	01:36:45.81	+15:47:40.2	26.12	1.73	>3.45	4.49
7	01:36:47.57	+15:47:15.1	26.13	>8.35	...	...
8 ^a	01:36:44.42	+15:45:15.6	26.17	6.86	>1.27	>2.16
9 ^a	01:36:43.35	+15:48:04.6	26.21	7.49	>0.73	>1.28
10	01:36:41.77	+15:47:11.6	26.23	4.83	>1.11	>1.94
11 ^a	01:36:43.33	+15:45:12.2	26.29	10.18	>0.77	>1.31
12	01:36:41.56	+15:47:18.1	26.30	>7.13	...	...
13	01:36:43.29	+15:47:57.3	26.32	1.17	>4.22	4.69
14 ^a	01:36:47.02	+15:46:57.5	26.36	>6.76	...	...
15 ^a	01:36:47.27	+15:47:39.7	26.37	>6.67	...	...
16	01:36:43.72	+15:47:38.7	26.38	>6.62	...	...
17	01:36:43.23	+15:47:58.5	26.39	>6.59	...	...
18	01:36:41.76	+15:47:13.1	26.50	>5.95	...	...
19	01:36:45.66	+15:47:07.7	26.52	2.07	>1.99	1.26
20	01:36:41.78	+15:47:15.9	26.52	3.72	>1.10	>1.93
21 ^a	01:36:45.34	+15:47:34.5	26.56	0.77	>5.15	3.69
22	01:36:43.50	+15:47:23.3	26.61	0.86	3.63	1.99
23 ^a	01:36:43.14	+15:47:40.3	26.66	>5.13	...	...
24	01:36:42.27	+15:47:27.6	26.73	>4.80	...	...
25	01:36:44.14	+15:47:08.7	26.77	>4.61	...	...
26	01:36:45.05	+15:46:53.9	26.77	>4.61	...	...
27	01:36:44.36	+15:45:06.5	26.81	>6.33	...	...
28	01:36:42.22	+15:47:47.5	26.81	3.47	>0.90	>1.58
29	01:36:42.89	+15:45:18.7	26.84	>6.14	...	...
30	01:36:43.89	+15:47:16.9	26.86	>4.27	...	...
31	01:36:44.58	+15:47:12.5	26.86	>4.26	...	...
32	01:36:41.95	+15:47:38.9	26.90	2.52	>1.15	>2.01
33	01:36:44.07	+15:47:06.5	26.91	2.23	>1.29	>2.27
34	01:36:45.12	+15:47:25.0	26.92	>4.04	...	...
35	01:36:41.29	+15:47:54.5	26.94	>3.94	...	...
36	01:36:42.86	+15:47:35.0	26.98	>3.82	...	...
37	01:36:46.79	+15:47:17.9	27.05	0.68	>3.68	5.90
38	01:36:42.27	+15:47:08.5	27.07	0.88	>2.81	4.41
39 ^a	01:36:43.25	+15:44:58.0	27.10	>4.85	...	...
40	01:36:43.10	+15:47:45.6	27.18	>3.17	...	...
41	01:36:45.24	+15:45:15.0	27.18	>4.51	...	...
42	01:36:43.79	+15:48:01.2	27.23	0.48	>4.45	2.39
43	01:36:47.61	+15:47:03.0	27.26	1.01	>2.06	2.32
44	01:36:39.68	+15:47:49.1	27.27	1.33	>1.55	1.92
45	01:36:40.35	+15:47:44.8	27.27	>2.91	...	...
46	01:36:41.38	+15:47:25.9	27.37	0.42	>4.51	5.12
47	01:36:42.88	+15:48:02.5	27.39	>2.60	...	...
48	01:36:45.58	+15:47:06.2	27.41	>2.57	...	...
49	01:36:40.46	+15:47:58.6	27.41	>2.57	...	...
50	01:36:41.33	+15:45:23.1	27.42	>3.61	...	...
51	01:36:41.42	+15:47:33.3	27.52	>2.32	...	...
52	01:36:42.56	+15:48:05.6	27.55	>2.25	...	...
53	01:36:43.21	+15:48:05.8	27.57	>2.22	...	...
54	01:36:42.59	+15:47:27.6	27.57	>2.21	...	...
55	01:36:41.26	+15:47:21.2	27.57	>2.21	...	...
56	01:36:42.81	+15:45:11.9	27.62	0.41	>5.59	4.92
57	01:36:42.79	+15:48:01.1	27.66	>2.05	...	...
58	01:36:42.67	+15:45:46.0	27.67	2.04	>1.08	1.00
59	01:36:46.71	+15:47:16.0	27.71	>1.95	...	...
60	01:36:42.65	+15:44:49.2	27.73	>2.71	...	...
61	01:36:43.63	+15:47:45.7	27.90	>1.63	...	...
62	01:36:42.68	+15:47:36.6	27.90	1.02	>1.12	1.09
63	01:36:41.85	+15:48:07.9	27.92	>1.60	...	...

Note.

^a Identified in Herrmann & Ciardullo (2009).

good (Figure 8), and we measure a distance modulus of $29.65^{+0.08}_{-0.13}$ mag, in very good agreement with Herrmann et al. (2008).

We also compare our measured distance modulus to results from the literature, using a variety of methods (Figure 10; Table 4). We find good agreement with the most recent SN II

Figure 8. Planetary nebulae luminosity function using only our 63 identified PNe (black), including contamination from 26 SNRs (red). Poisson error bars are included for each bin. Our fits (dotted lines) are done using a maximum likelihood method on our sources brighter than the 27 mag completeness limit (filled symbols), allowing us to measure a distance modulus of $29.91^{+0.08}_{-0.13}$ mag for NGC 628. Introducing the SNR contaminants, we find a distance modulus of $29.65^{+0.08}_{-0.13}$ mag, reproducing the previous PNLF estimate based on narrowband imaging (Herrmann et al. 2008).

Figure 9. Cumulative PNe luminosity function for our 63 identified PNe (black) and including contamination from 26 SNRs (red). Although our low number statistics make it difficult to judge the quality of the PNLF fit in Figure 8, we see here the good agreement between our fit (dotted lines) and the observed populations, down to our completeness limit of 27 mag.

methods, and 2σ agreement with the recent TRGB measurement, which has much smaller uncertainties than all other methods.

4.2. An IFU Approach to PN Studies

Instruments like MUSE, with high spatial resolution and large field of view, are well suited to identifying PNe populations within nearby galaxies, but there has been limited application of such an IFU approach so far (Roth et al. 2004; Sarzi et al. 2011). We have shown that, at a physical resolution of 50 pc, this technique can recover the PNLF cleaned of

Figure 10. Comparison of our distance estimate (filled circle) to values in the literature. We show 1σ uncertainties for each measurement. Our revised PNLF distance is in much closer agreement (within 2σ) with the tip of the red giant branch (TRGB) measurement (triangle) than the previous PNLF distance (open circle). We also find good agreement with most of the distances measured using the brightest supergiant (BS, squares) and SN II (diamonds) methods, although these all have much larger error bars. See Table 4 for detailed references for all methods.

contaminants to measure a robust distance to NGC 628. This is most important for galaxies with a substantial amount of cold ISM, where supernova ejecta shocks are likely to create [O III] 5007 emission. In systems without a cold ISM (elliptical galaxies, spiral bulges), the ejecta expands more freely, and will not normally create [O III] emission. Given the large field covered by our observations, we explore at what spatial scale our simple method for selecting PNe (as outlined in this paper) breaks down.

At 100 pc scales ($2''$ resolution), we can identify 36 [O III] sources with line ratios consistent with PNe. In most cases, both the $H\alpha$ and [S II] lines are not detected above the background diffuse emission, making removal of SNR contaminants difficult without more detailed background modeling. We find that 23 of these objects are in our PNe catalog, and four are in our SNR catalog. The remaining are typically within $\sim 5''$ from a cataloged PN, and are likely PNe that are blended with a neighboring object, shifting the centroid.

At 150 pc scales ($3''$ resolution), we identify 24 [O III] sources with line ratios consistent with PNe, 13 of which are in our PNe catalog and three of which are in the SNR catalog.

At scales larger than 150 pc, we see extensive blending of the [O III] emission because we are tracing mainly the crowded, central, and highly star-forming regions. Crowded-field 3D spectrophotometry (Roth et al. 2004) can be used to remove background due to unresolved stars and diffuse ISM emission, and recover faint emission lines for detailed PN modeling. More involved source extraction techniques, including use of prior knowledge of source positions from other high-resolution imaging (Kamann et al. 2013), would be necessary to extract useful PN line fluxes on scales larger than this.

Table 4
Comparison of Distance Moduli

$(m - M)$ (mag)	D (Mpc)	Method	Reference
29.32 ± 0.40	7.31 ± 1.35	Brightest supergiants	Sharina et al. (1996)
29.44 ± 0.48	7.73 ± 1.71	Brightest supergiants	Hendry et al. (2005)
29.91 ± 0.49	9.59 ± 2.16	Brightest supergiants (alternate calibration)	Hendry et al. (2005)
29.91 ± 0.63	9.59 ± 2.78	SN II standard candle	Hendry et al. (2005)
$29.67^{+0.06}_{-0.07}$	$8.59^{+0.24}_{-0.28}$	PNLF	Herrmann et al. (2008)
29.98 ± 0.28	9.91 ± 1.28	SN II standard candle	Olivares et al. (2010)
29.76 ± 0.24	8.95 ± 0.99	SN II photospheric magnitude	Rodríguez et al. (2014)
29.74 ± 0.18	8.87 ± 0.74	SN II photospheric magnitude	Rodríguez et al. (2014)
30.01 ± 0.07	10.05 ± 0.32	SN II photospheric magnitude	Rodríguez et al. (2014)
29.88 ± 0.05	9.46 ± 0.22	SN II photospheric magnitude	Rodríguez et al. (2014)
30.04 ± 0.03	10.19 ± 0.14	TRGB	Jang & Lee (2014)
$29.91^{+0.08}_{-0.13}$	$9.59^{+0.35}_{-0.57}$	PNLF	this work

As the PN population traces the stellar density, there is a clear advantage to pursuing PN studies within galaxy disks. With narrowband imaging, these environments were challenging due to the high stellar and ISM background. However, with full spectroscopic information the search for PNe in these regions can recover a large number of objects with relatively small spatial coverage. For our MUSE pointings, we measure the stellar mass surface density based on the Spitzer imaging where emission from dust and other nonstellar sources has been removed (Querejeta et al. 2015), and we estimate the V-band surface brightness based on a published radial profile (Fisher & Drory 2008). Our southern pointing (at slightly larger radius) has an average stellar surface density of $6.9 \times 10^7 M_{\odot} \text{ kpc}^{-2}$ ($22.5 \text{ mag arcsec}^{-2}$) and contains 1.6 PNe per kpc^2 , whereas the inner pointings at roughly twice the stellar surface density ($1.9 \times 10^8 M_{\odot} \text{ kpc}^{-2}$, $21.0 \text{ mag arcsec}^{-2}$) contain twice the PN density (3.1 PNe per kpc^2). Counting only the PNe detected above our completeness limit and suitable for fitting the PNLf, we find an average of 1.3 PNe per kpc^2 in the inner pointings.

For a galaxy twice as distant as NGC 628 ($\sim 20 \text{ Mpc}$), the PNLf cutoff would be 1.5 mag fainter. To robustly trace the PNLf cutoff the required PN survey depth is 0.8 mag, however, even a small sample (20 PNe) can still accurately recover the distance to within 3% with only $\sim 10\%$ increased uncertainty (Jacoby 1997, pp. 197–211). Given our observed PN density, this corresponds to a survey area of $\sim 15 \text{ kpc}^2$ within the central galaxy disk. For our shallower observations (42 minutes on source) with 27.0 mag completeness limit, this means we could measure a distance out to $D = 13.6 \text{ Mpc}$ ($1'' = 80 \text{ pc}$). Our deeper southern pointings (50 minutes on source), reaching 27.5 mag completeness limit, could measure a distance out to $D = 17.1 \text{ Mpc}$ ($1'' = 80 \text{ pc}$). In both cases, a single MUSE observation in Wide Field Mode would cover the required survey area at a spatial resolution ($< 100 \text{ pc}$) sufficient for clean separation of a significant number of PNe.

Optical IFU identification of PNe has applications, not just for measuring distances, but also for using PNe to trace chemical enrichment through detailed study of their metallicity (Ali et al. 2016, and references therein), and stellar disk dynamics through study of the emission line kinematics (Cortesi et al. 2013, and references therein).

5. CONCLUSION

We use the optical IFU instrument MUSE on the VLT to map [O III], H α , [N II] and [S II] emission lines across 27 kpc^2

within the central star-forming galaxy disk of NGC 628. We select unresolved [O III] emission line sources, and distinguish 87 H II regions and 30 SNRs (Table 2) as contaminants to our sample, using diagnostic line ratios.

In total, we identify 63 PNe within our fields (Table 3), and we determine a new PNLf distance modulus of $29.91^{+0.08}_{-0.13} \text{ mag}$ ($9.59^{+0.35}_{-0.57} \text{ Mpc}$) by using the 36 PNe with $m_{[\text{O III}]}$ $< 27.0 \text{ mag}$. This is in good agreement with literature values, particularly with the TRGB estimate (which has the smallest uncertainty). Our measured distance is significantly larger than the previously reported PNLf distance; however, we can reproduce this lower distance estimate by biasing our sample with SNR contaminants, demonstrating the power of full spectral information over narrowband imaging when isolating the PN sample.

Given our limited spatial coverage within the Galaxy, we show that this technique can be used to refine distance estimates, even when only limited IFU observations are available. In particular, when targeting central regions where the stellar density is high, IFU techniques can recover large numbers of PNe ($\sim 3 \text{ PNe per kpc}^2$) within galaxies at 10–20 Mpc distances (50–100 pc at $1''$ resolution) that were previously unobservable due to confusion with background stellar and ISM emission.

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